

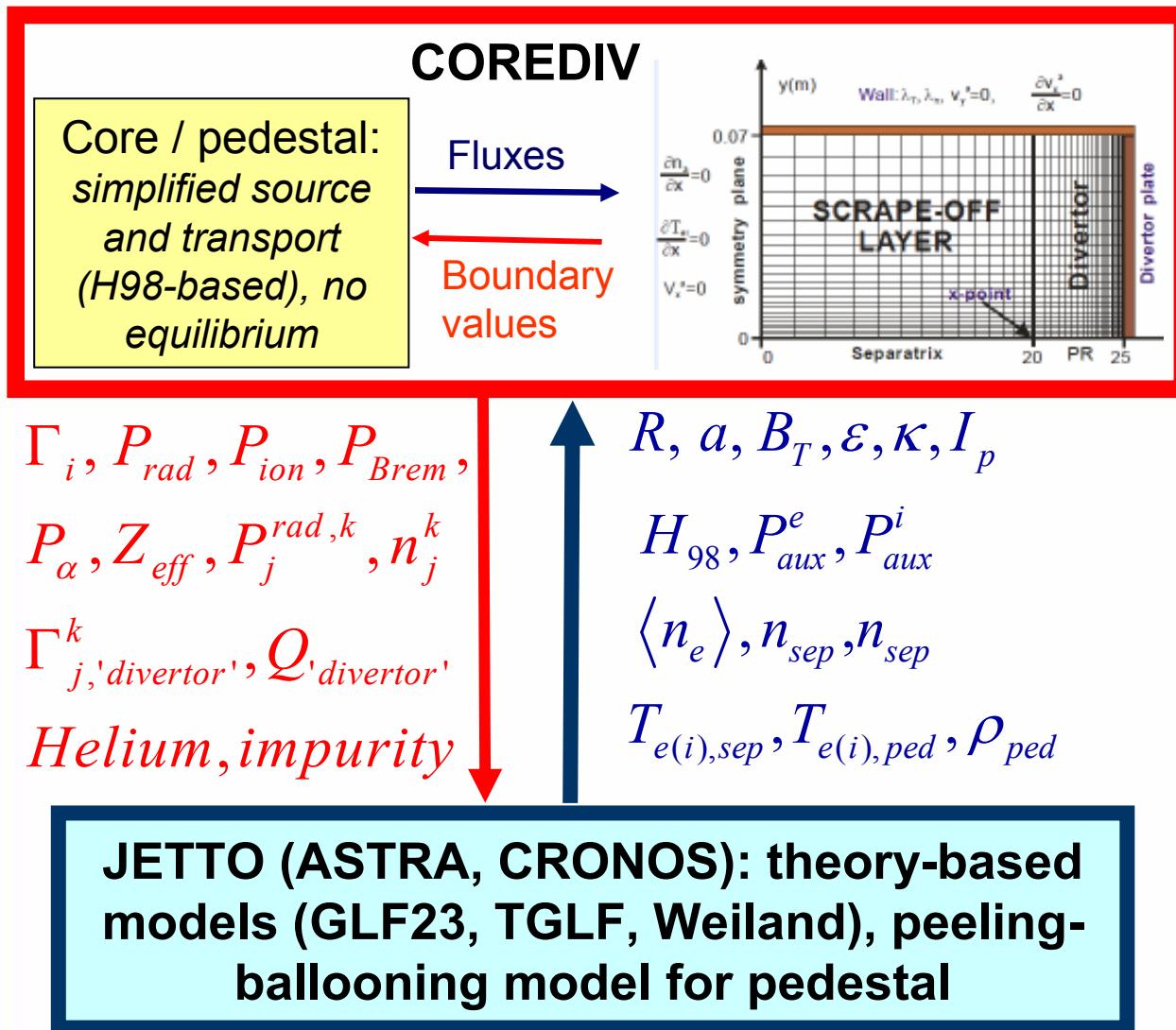
Integrated core-SOL modelling including impurity: ITER H-mode plasma

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Approaches to integrated core-SOL modelling

- **Coupling of complex core-SOL codes (eg. JINTRAC):**
 - *self-consistent simulations of strongly non-linearly coupled core transport, H&CD, pedestal and SOL physics is an advantage*
 - *2D SOL simulations are computationally expensive. Possible issues with core-SOL convergence*
 - *limited application, mainly for transients: ELMs, L-H transition, ...*
- **One of alternative approaches: scaling for SOL integrated in core codes**
- **Combination of core code (ASTRA, CRONOS, JETTO) and COREDIV as a compromise: sufficient physics complexity and acceptable simulation speed**
 - *core transport physics: theory-based transport, accurate H&CD simulations, equilibrium, current diffusion*
 - *impurity equations (COREDIV)*
 - *SOL-divertor (COREDIV)*
- **Application to ITER H-mode (long pulse scenario): effect of pedestal density on plasma performance**

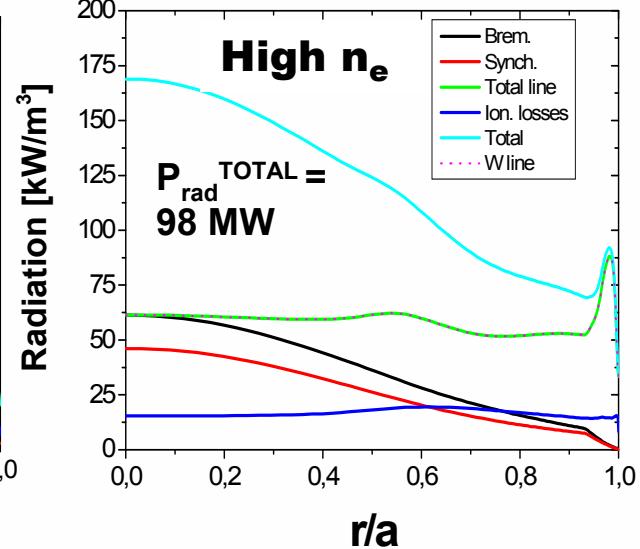
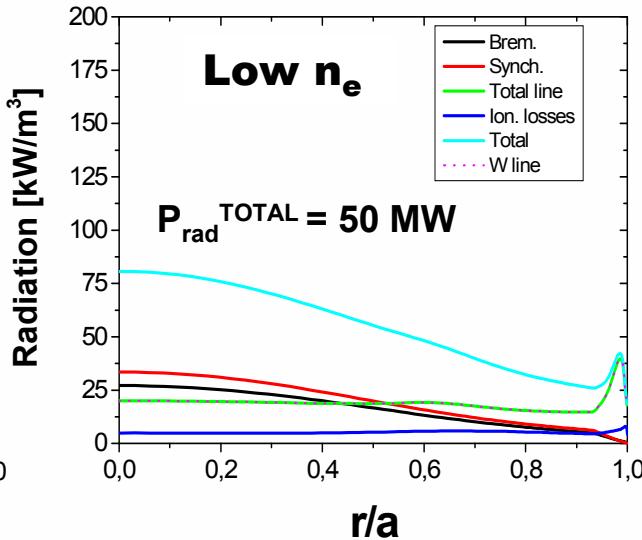
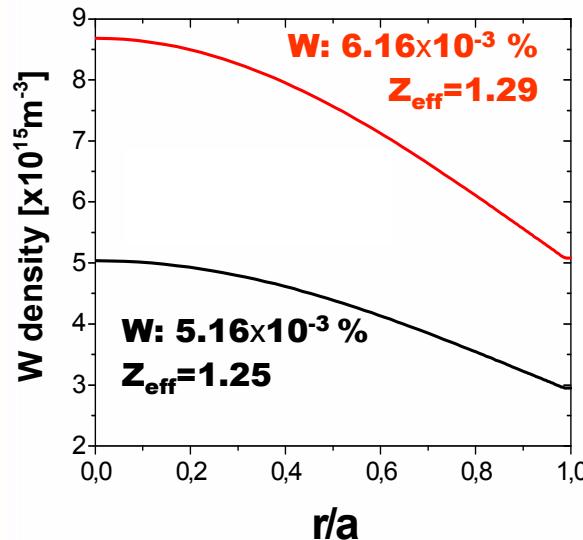
Coupled JETTO-COREDIV simulations for ITER H-mode plasmas



- **ITER H-mode scenario:** 15 MA, 5.34 T, 33 MW (NBI) + 20 MW (ECRH)
- **Low density (ref.):** $n_{e,ped} = 6e19 \text{ m}^{-3}$, $T_{e,ped} = 6.7 \text{ keV}$, $T_{i,ped} = 7.5 \text{ keV}$
- **High density:** $n_{e,ped} = 9.5e19 \text{ m}^{-3}$, $T_{e,ped} = 4.35 \text{ keV}$, $T_{i,ped} = 4.83 \text{ keV}$
- **JETTO / GLF23, EPED, SOLPS boundary, no W**
 $\rightarrow H98=0.85$ at low and high ne, Q =
- **COREDIV: effect of W?**

ITER H-mode: COREDIV simulations

I. Ivanova-Stanik, F. Köchl et al, PET 2013



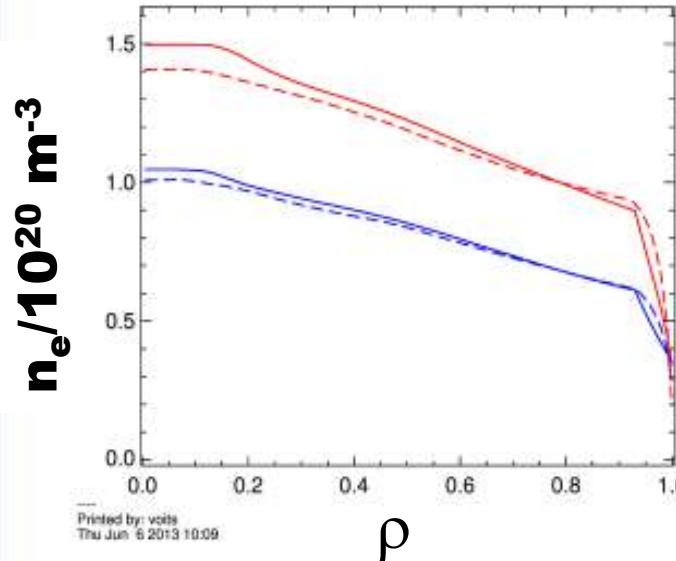
- Higher density is achieved by increasing gas puff (feedback: $(8.8 \rightarrow 21.2) \times 10^{21} \text{ 1/s}$)
- $D_{D,T} = 0.35 \chi_{i,e}$; $V_{D,T}$ adjusted to match the GLF23 n_e peaking. Similar particle transport at low and high n_e . Impurity transport is the same as main species
- $P_\alpha \uparrow$, power to plate \uparrow , W sputtering \uparrow with $n_e \rightarrow$ high n_W and P_{rad} . W concentration increases due to sputtering, not transport
- Increase of power to plate $\sim 11\%$ (strongly increased α -heating and W radiation partly compensate each other)

	Low n_e	High n_e
$P_\alpha [\text{MW}]$	79.4	140.1
Core radiation [MW]	41.92	89.97
Radiation fraction	0.378	0.505
Power to plate [MW]	76.05	85.32
T_e plate [eV]	25.45	33.55
n_e plate [$\times 10^{20} \text{ m}^{-3}$]	2.1	1.43
Recycling coefficient	0.9958	0.989
Sputtering flux, 1/s	1.1×10^{20}	1.9×10^{20}

Work in progress

JETTO simulations with W radiation from COREDIV (F. Koechl)

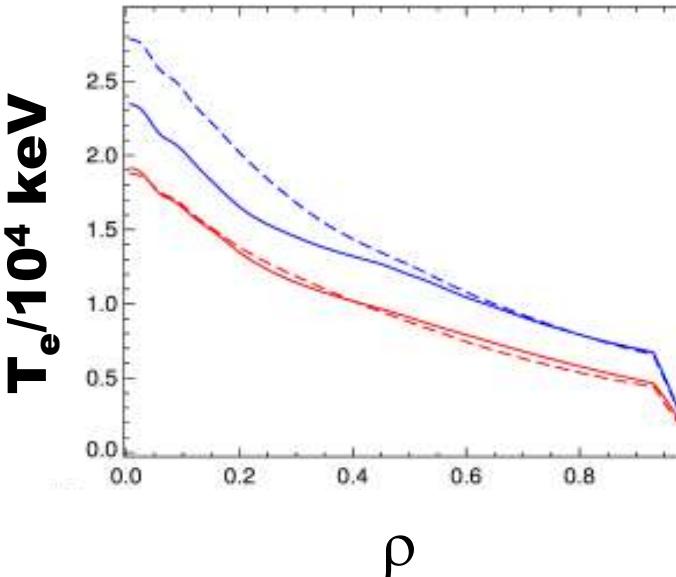
JET Data Display



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03298 JSP/NE
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T=481,407
03298 JSP/NE
Seq=19445 (0) Uld=Kochl
T=484,824

$n_e / 10^{20} \text{ m}^{-3}$

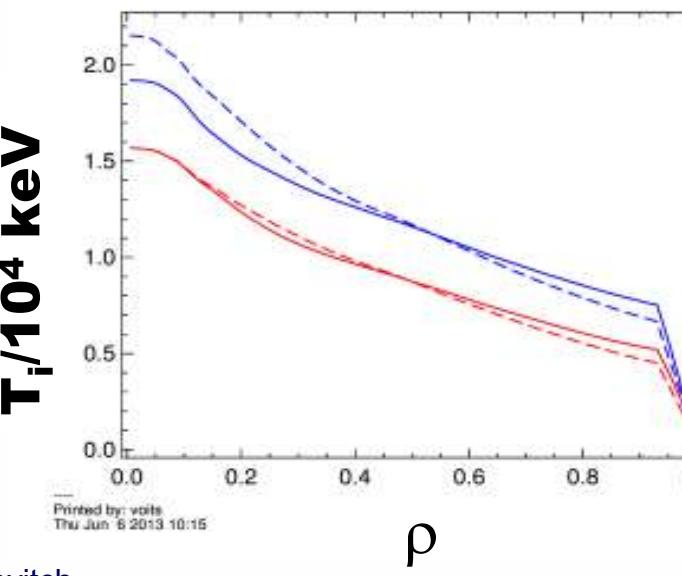
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03298 JSP/TE
Seq=19446 (0) Uld=Kochl
T=481,407
03298 JSP/TE
Seq=19445 (0) Uld=Kochl
T=484,824

$T_e / 10^4 \text{ keV}$

ρ



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Sensitivity to local transport (impurity, main species and thermal) is particularly high with dominant alpha heating

- **High sensitivity to $D_{D,T}/\chi_{i,e}$. In present ITER simulations $D_{D,T}/\chi_{i,e} = 0.35$**
 - *when $D_{D,T}/\chi_{i,e} = 0.1$ no self-consistent solution at high density (large He concentration, low P_α , high radiation)*
 - *larger $D_{D,T}/\chi_{i,e}$ margins for flat ne profiles → transport of main species determining the ne peaking is important*
 - *difficult to reduce $D_{D,T}/\chi_{i,e}$ uncertainty using present experiments ($D_{D,T}/\chi_{i,e}=0.1-0.35$ gives similar results for JET plasmas)*
- **Fusion performance is sensitive to core thermal transport model: at given H98 factor the GLF23 and scaling-based models predict different Ti shape and P_α ($P_\alpha=84$ MW (GLF23) and $P_\alpha=140$ MW (scaling-based))**